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LETTER TO THE EDITOR

The sine-Gordon equation and the trace method

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**Abstract.** The trace method which has been proposed by Wadati and Sawada is applied to the sine-Gordon equation. The  $N$  soliton solution is derived.

The trace method which has been applied to the  $\kappa\Delta V$  [1],  $M\kappa\Delta V$  [2],  $\kappa P$  [3] and some other equations [4] is useful for understanding these equations [1]. The  $N$  soliton solution and some other results of these equations have been derived through the trace method.

The present letter deals with an application of the trace method to the sine-Gordon equation as follows:

$$U_{xt} + \sin U = 0 \quad U(x, t) \rightarrow \text{constant as } x \rightarrow \infty. \tag{1}$$

Wadati and Sawada have pointed out that it is very difficult to solve the sine-Gordon equation by means of the trace method directly [2]. In order to avoid dealing with the transcendental non-linearity  $\sin U$ , it is convenient to re-express the equation in terms of  $V = U_x$ ,  $W = \cos U$ , i.e.

$$V_{xt} = -WV \quad W_x = V_t V. \tag{2}$$

Substituting the formal series

$$\begin{aligned} V &= V^{(1)} + V^{(3)} + \dots + V^{(2n+1)} + \dots \\ W &= 1 + W^{(2)} + V^{(4)} + \dots + W^{(2n)} + \dots \end{aligned} \tag{3}$$

into equations (2), we obtain a set of equations for  $W^{(2n)}$  and  $V^{(2n+1)}$  ( $n = 0, 1, 2, \dots$ ):

$$\begin{aligned} W^{(0)} &= 1 & V_{xt}^{(1)} &= -V^{(1)} & W_x^{(2)} &= V_t^{(1)} V^{(1)} \\ V_{xt}^{(3)} &= -V^{(3)} - W^{(2)} V^{(1)} \\ &\vdots \\ W_x^{(2n)} &= \sum_{r=0}^{n-1} V_t^{(2r+1)} V^{(2n-2r-1)} \end{aligned} \tag{4}$$

$$\begin{aligned} V_{xt}^{(2n+1)} &= -\sum_{r=0}^n W^{(2r)} V^{(2n-2r+1)} \\ &\vdots \end{aligned} \tag{5}$$

We can solve the set of equations iteratively:

$$\begin{aligned}
 V^{(1)} &= 4 \sum_{r=1}^N \Phi_r^2(x, t) \\
 W^{(2)} &= - \sum_{m=1}^N \sum_{n=1}^N \frac{\Phi_m^2(x, t)\Phi_n^2(x, t)}{(k_m + k_n)} (k_m^{-1} + k_n^{-1}) \\
 V^{(3)} &= -4 \sum_{r=1}^N \sum_{m=1}^N \sum_{n=1}^N \frac{\Phi_r^2(x, t)\Phi_m^2(x, t)\Phi_n^2(x, t)}{(k_r + k_m)(k_m + k_n)}.
 \end{aligned}$$

Here  $\Phi_n = A_n(0) \exp(k_n x - t/4k_n) = \Phi_n(x, t)$ ,  $k_n$  are positive and distinct constants  $A_n(0)$  are real or pure imaginary ( $n = 1, 2, \dots, N$ ). We introduce an  $N \times N$  matrix whose elements are given by  $B_{mn} = \int_{-\infty}^x \Phi_m(y, t)\Phi_n(y, t) dy = [1/(k_m + k_n)]\Phi_m(x, t)\Phi_n(x, t)$ .

With the matrix  $B$ ,  $V^{(1)}$ ,  $W^{(2)}$ , and  $V^{(3)}$  are expressed as

$$V^{(1)} = 4 \text{Tr}(B_x) \qquad W^{(2)} = 2\text{Tr}(B_{xt}^2) \qquad V^{(3)} = -4\text{Tr}(B_x B^2).$$

In general, we can prove that

$$V^{(2n+1)} = 4(-1)^n \text{Tr}(B_x B^{2n}) \tag{6}$$

$$W^{(2n)} = (2/n)(-1)^{n-1} \text{Tr}(B_{xt}^{2n}) \tag{7}$$

satisfy equations (4) and (5). First, it is easy to see that

$$\text{Tr}[(B^{2n+1})_x] = (2n + 1) \text{Tr}(B_x B^{2n}). \tag{8}$$

Therefore

$$\begin{aligned}
 V^{(2n+1)} &= [4/(2n + 1)](-1)^n \text{Tr}[(B^{2n+1})_x] \\
 V_{xt}^{(2n+1)} &= [4/(2n + 1)](-1)^n [\text{Tr}(B^{2n+1})]_{txx}.
 \end{aligned} \tag{9}$$

Now we make the following definitions.

- (i)  $\sigma$  is a cyclic transformation,  $\sigma(1, 2, \dots, 2n + 1) = (2, 3, \dots, 2n + 1, 1)$ .
- (ii)  $T$  is an operator for a subscript variable, defined by  $T(F(k_{r_1}, k_{r_2}, \dots, k_{r_{2n+1}})) = F(k_{r_{\sigma(1)}}, k_{r_{\sigma(2)}}, \dots, k_{r_{\sigma(2n+1)}})$ .
- (iii)  $S$  is an order symmetry operator defined by

$$S = T^1 + T^2 + \dots + T^{2n+1}.$$

Obviously, we have that

- (1)  $S$  is a linear operator
- (2)  $ST \neq S$ ,  $TS = S$  and
- (3) if  $f$  is a function of the  $(2n + 1)$ th subscript variable and  $T(f) = f$ , then  $S(f) = (2n + 1)f$  and  $S(gf) = (S(g))f$ , where  $g$  is a function of the subscript variable.

From (8), we have

$$W^{(2n)} = (2/n)(-1)^{n-1} \text{Tr}(B_{xt}^{2n}) = 4(-1)^{n-1} [\text{Tr}(B_x B^{2n-1})]_t. \tag{10}$$

In the following we calculate the right-hand side of equation (5) and simplify the expression by writing  $1, 2, \dots, 2n+1$  instead of  $r_1, r_2, \dots, r_{2n+1}$ :

$$\begin{aligned}
 & - \sum_{r=0}^n W^{(2r)} V^{(2n-2r+1)} \\
 &= 4(-1)^{n+1} \sum_1 \sum_2 \dots \sum_{2n+1} \\
 & \quad \times \left( 1 + 2 \sum_{r=1}^n (k_1^{-1} + k_2^{-1} + \dots + k_{2r}^{-1})(k_{2r} + k_{2r+1}) \right) \Phi_1^2 \Phi_2^2 \dots \Phi_{2n+1}^2 \\
 & \quad \times [(k_1 + k_2)(k_2 + k_3) \dots (k_{2n} + k_{2n-1})]^{-1} \\
 &= 94(-1)^{n+1} \sum_1 \sum_2 \dots \sum_{2n+1} \\
 & \quad \times \left( 1 + 2 \sum_{r=1}^n (k_1^{-1} + k_2^{-1} + \dots + k_{2r}^{-1})(k_{2r} + k_{2r+1}) \right) \\
 & \quad \times (k_{2n+1} + k_1) \Phi_1^2 \Phi_2^2 \dots \Phi_{2n+1}^2 \\
 & \quad \times [(2n+1)(k_1 + k_2)(k_2 + k_3) \dots (k_{2n} + k_{2n+1})(k_{2n+1} + k_1)]^{-1}.
 \end{aligned}$$

From the well known Abelian transformation, we have

$$\begin{aligned}
 \sum_{r=1}^n (k_1^{-1} + k_2^{-1} + \dots + k_{2r}^{-1})(k_{2r} + k_{2r+1}) &= \sum_{i=1}^n (k_{2i-1}^{-1} + k_{2i}^{-1}) \left( \sum_{r=i}^n (k_{2r} + k_{2r+1}) \right) \\
 S \left( 1 + 2 \sum_{r=1}^n (k_1^{-1} + k_2^{-1} + \dots + k_{2r}^{-1})(k_{2r} + k_{2r+1}) \right) (k_{2n+1} + k_1) \\
 &= S(k_{2n+1} + k_1) + 2S(k_{2n+1} + k_1) \sum_{i=1}^n (k_{2i-1}^{-1} + k_{2i}^{-1}) \left( \sum_{r=i}^n (k_{2r} + k_{2r+1}) \right) \\
 &= 2S(k_1) + 2S \sum_{i=1}^n \left[ T^{2(n-i)+3}(k_{2n+1} + k_1) k_{2i-1}^{-1} \left( \sum_{r=i}^n (k_{2r} + k_{2r+1}) \right) \right. \\
 & \quad \left. + T^{2(n-i+1)}(k_{2n+1} + k_1) k_{2i} \left( \sum_{r=i}^n (k_{2r} + k_{2r+1}) \right) \right] \\
 &= 2S \left( k_1 + k_1^{-1} \sum_{i=0}^{n-1} \sum_{r=1}^{n-i} [(k_{2r} + k_{2r+1})(k_{2n-2i+1} + k_{2n-2i+2}) \right. \\
 & \quad \left. + (k_{2r-1} + k_{2r})(k_{2n-2i} + k_{2n-2i+1})] \right) \\
 &= 2Sk_1 \left( k_1 + \sum_{i=0}^{n-1} \sum_{r=2}^{2n-2i+1} k_r k_{2n-2i+1} + \sum_{i=1}^{n-1} \sum_{r=2}^{2n-2i+1} k_r k_{2n-2i+2} \right. \\
 & \quad \left. + \sum_{r=2}^{2n+1} k_r k_1 + \sum_{i=0}^{n-1} \sum_{r=1}^{2n-2i} k_r k_{2n-2i} + \sum_{i=0}^{n-1} \sum_{r=1}^{2n-2i} k_r k_{2n-2i+1} \right) \\
 &= 2Sk_1^{-1} \left( \sum_{i \leq j} k_i k_j + \sum_{i > j} k_i k_j \right) \\
 &= 2(k_1^{-1} + k_2^{-1} + \dots + k_{2n+1}^{-1})(k_1 + k_2 + \dots + k_{2n+1})^2.
 \end{aligned}$$

Therefore, we arrive at

$$\begin{aligned}
 - \sum_{r=0}^n W^{(2r)} V^{(2n-2r+1)} & \\
 &= [4/(2n+1)](-1)^n \text{Tr}(B_{xx}^{2n+1}) \\
 &= V_{xt}^{(2n+1)}.
 \end{aligned}$$

Next we shall show that (6) and (7) satisfy equation (4):

$$\begin{aligned}
 \sum_{r=0}^{n-1} V_t^{(2r+1)} V^{(2n-2r-1)} & \\
 &= \frac{1}{2} \left( \sum_{r=0}^{n-1} V_t^{(2r+1)} V^{(2n-2r-1)} + \sum_{r=0}^{n-1} V^{(2r+1)} V_t^{(2n-2r-1)} \right) \\
 &= 8(-1)^{n-1} \sum_{r=0}^{n-1} \left( \sum_1 \sum_2 \dots \sum_{2n} [(k_1+k_2) \dots (k_{2r}+k_{2r+1}) \right. \\
 &\quad \left. \times (k_{2r+2}+k_{2r+3}) \dots (k_{2n-1}+k_{2n})]^{-1} \Phi_1^2 \Phi_2^2 \dots \Phi_{2n}^2 \right)_t \\
 &= 8(-1)^{n-1} \left( \sum_1 \sum_2 \dots \sum_{2n} \sum_{r=0}^{n-1} (k_{2r+1}+k_{2r+2}) [(k_1+k_2)(k_2+k_3) \dots \right. \\
 &\quad \left. \times (k_{2n-1}+k_{2n})]^{-1} \Phi_1^2 \Phi_2^2 \dots \Phi_{2n}^2 \right)_t \\
 &= 4(-1)^{n-1} [\text{Tr}(B_x B^{2n-1})]_{xt} \\
 &= W_x^{(2n)}.
 \end{aligned}$$

We have proved that (6) and (7) satisfy equations (4) and (5). Therefore we obtain the solution for equations (2) in the following form:

$$\begin{aligned}
 V &= 4 \text{Tr}(B_x - B_x B^2 + B_x B^4 + \dots) \neq 4 \text{Tr}[B_x(1+B^2)^{-1}] \\
 &= 4[\text{Tr}(\tan^{-1} B)]_x = \frac{2}{i} \left[ \text{Tr} \log \left( \frac{I+iB}{I-iB} \right) \right]_x
 \end{aligned}$$

$$W = 1 + 2\{\text{Tr}[B^2 - \frac{1}{2}B^4 + \dots + (-1)^{n-1}n^{-1}B^{2n} + \dots]\}_{xt} = 1 + 2[\text{Tr} \log(1+B^2)]_{xt}.$$

From  $V = U_x$ , we have

$$U = 4 \text{Tr}(\tan^{-1} B) = \frac{2}{i} \text{Tr} \log \left( \frac{I+iB}{I-iB} \right). \tag{11}$$

We have to show that the above  $U$  satisfy equation (1). Due to  $V = U_x$ , we have

$$W_x = V_t V = U_{xt} U_x \tag{12}$$

$$W_{xx} = U_{xxt} U_x + U_{xt} U_{xx}.$$

Using  $V_{xt} = -WV$ , i.e.  $U_{xxt} = -WU_x$  we obtain

$$W_{xx} = -WU_x^2 + U_{xt} U_{xx}.$$

Therefore we have

$$\begin{aligned} 0 &= W_{xx}/U_x^2 + W - U_{xt}U_{xx}/U_x^2 = W_{xx}/U_x^2 - W_xU_{xx}/U_x^3 + W \\ &= (W_x/U_x)_x \frac{dX}{dU} + W = \frac{d^2W}{dU^2} + W. \end{aligned}$$

We solve  $W_{uu} + W = 0$  and obtain

$$W = C_1 \cos U + C_2 \sin U \quad C_1 \text{ and } C_2 \text{ constant.}$$

From (11) we see that  $U$  is an odd function of  $B$ , but  $W$  is an even function of  $B$ . Hence  $W$  is an even function of  $U$  and  $C_2 = 0$ . If we let  $U = 0$ , then from (11)  $B = 0$ , and  $C_1 = C_1 \cos 0 = 1 = W$  (as  $B = 0$ ). Therefore we arrive at  $W = \cos U$ . Using  $W_x = V_t V$ , we have  $(\cos U)_x = U_{xt}U_x = (-\sin U)U_x$ , i.e.  $U_{xt} + \sin U = 0$ . In conclusion, the  $U$  defined by (11) satisfy equation (1).

We can easily derive the Gelfand-Levitan integral equation and the eigenvalue problem associated with the sine-Gordon equation in a method similar to the  $m\kappa\lambda v$  equation, by replacing  $U$  by  $V/2$  [2].

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## References

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